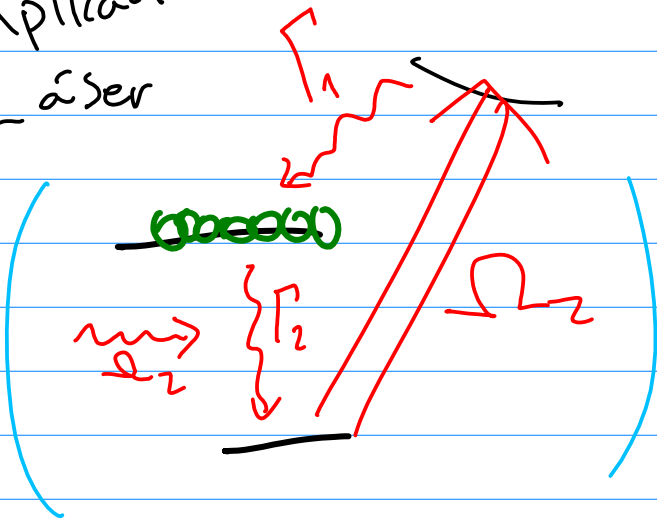


Átomo de 3 niveles

Aplicación.

Láser

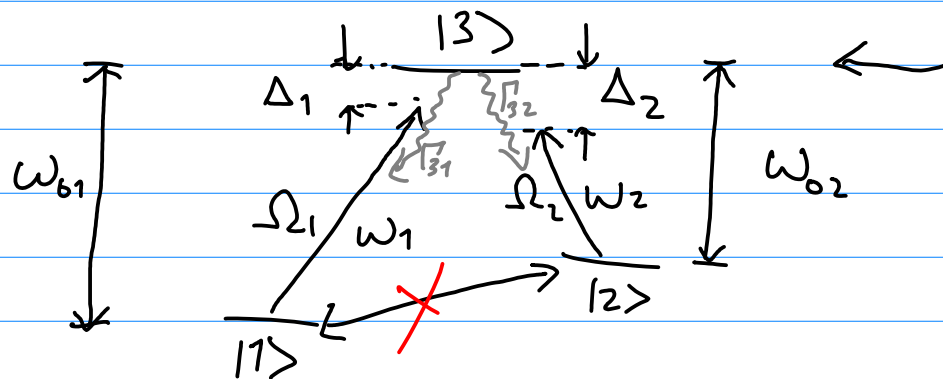
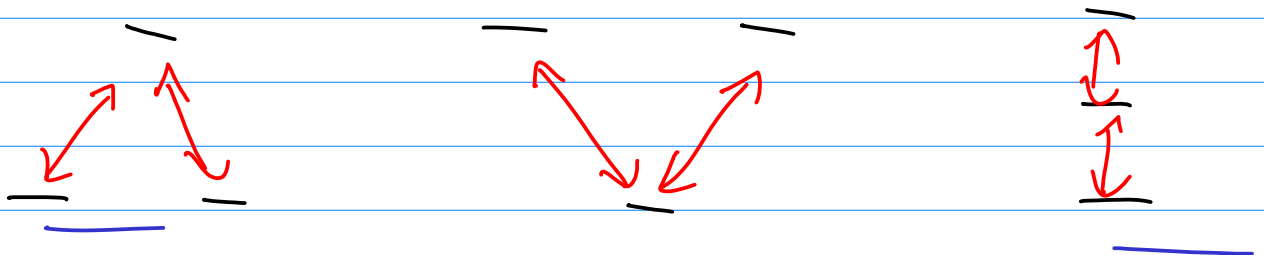


En átomo de 2 niveles.

$$\rho_{ee}(t \rightarrow \infty) = \frac{\Omega^2/\Gamma^2}{1 + (\frac{\Omega_0}{\Gamma})^2 + 2(\frac{\Omega}{\Gamma})^2}$$

Cuando $\Omega \rightarrow \infty$ $\rho_{ee} \rightarrow \frac{1}{2}$

Hay 3 tipos: Δ V I



$$\vec{E}(\vec{r}, t) = \hat{e}_1 \mathcal{E}_{o1} \cos(\vec{k}_1 \cdot \vec{r} - \omega_1 t) + \hat{e}_2 \mathcal{E}_{o2} \cos(\vec{k}_2 \cdot \vec{r} - \omega_2 t)$$

$$= \mathcal{E}^{(+)}(\vec{r}, t) + \mathcal{E}^{(-)}(\vec{r}, t)$$

$$\vec{E}^{(\pm)}(\vec{r}, t) = \frac{1}{2} \left(\hat{e}_1 \mathcal{E}_{o1} e^{\pm i \vec{k}_1 \cdot \vec{r}} e^{\mp i \omega_1 t} + \hat{e}_2 \mathcal{E}_{o2} e^{\pm i \vec{k}_2 \cdot \vec{r}} e^{\mp i \omega_2 t} \right)$$

$$H_A = -\hbar \omega_{01} |1\rangle\langle 1| - \hbar \omega_{02} |2\rangle\langle 2|$$

$$H_{AF} = -\vec{d} \cdot \vec{E} \stackrel{\text{RWA}}{\approx} -\vec{d}^{(+)} \cdot \vec{E}^{(-)} - \vec{d}^{(-)} \cdot \vec{E}^{(+)}$$

$$\vec{d} = \mathbb{1} \vec{d} \mathbb{1} = (|1\rangle\langle 1| + |2\rangle\langle 2| + |3\rangle\langle 3|) \vec{d} (|1\rangle\langle 1| + |2\rangle\langle 2| + |3\rangle\langle 3|)$$

$$= \langle 1|\vec{d}|3\rangle |1\rangle\langle 3| + \langle 2|\vec{d}|3\rangle |2\rangle\langle 3| + \langle 3|\vec{d}|1\rangle |3\rangle\langle 1| + \langle 3|\vec{d}|2\rangle |3\rangle\langle 2|$$

\uparrow σ_{13} \uparrow σ_{23} \uparrow σ_{13}^+ \uparrow σ_{23}^+

$$\Omega_\alpha = -\frac{\langle \alpha | \hat{E}_\alpha \cdot \vec{d} | \beta \rangle}{\hbar} \epsilon_{\alpha\beta} \quad \alpha = 1, 2$$

$$H_{AF} \stackrel{\text{RWA}}{\approx} \frac{\hbar \Omega_1}{2} \left(\sigma_{13} e^{-i\vec{k}_1 \cdot \vec{r}} e^{i\omega_1 t} + \sigma_{13}^+ e^{i\vec{k}_1 \cdot \vec{r}} e^{-i\omega_1 t} \right) + \frac{\hbar \Omega_2}{2} \left(\sigma_{23} e^{-i\vec{k}_2 \cdot \vec{r}} e^{i\omega_2 t} + \sigma_{23}^+ e^{i\vec{k}_2 \cdot \vec{r}} e^{-i\omega_2 t} \right)$$

Cambio al marco rotante

Para el átomo de dos niveles usamos

$$U = e^{i\omega t |e\rangle\langle e|} = \begin{pmatrix} 1 & 0 \\ 0 & e^{i\omega t} \end{pmatrix}$$

$$\tilde{H} = U^\dagger H U + i(\partial_t U) U^\dagger$$

$$\hbar \begin{pmatrix} 0 & \frac{\Omega}{2} e^{i\omega t} \\ -\frac{\Omega}{2} e^{-i\omega t} & \omega_0 \end{pmatrix} \longrightarrow \hbar \begin{pmatrix} 0 & \frac{\Omega}{2} \\ \frac{\Omega}{2} & \omega_0 - \omega \end{pmatrix}$$

Para 3 niveles

$$U = \begin{matrix} & \begin{matrix} |1\rangle & |2\rangle & |3\rangle \end{matrix} \\ \begin{matrix} \langle 1| \\ \langle 2| \\ \langle 3| \end{matrix} & \begin{pmatrix} e^{-i\omega_1 t} & 0 & 0 \\ 0 & e^{-i\omega_2 t} & 0 \\ 0 & 0 & 1 \end{pmatrix} \end{matrix}$$

Calculando $\tilde{H} = U^\dagger H U + i(\partial_t U) U^\dagger$

$$\tilde{H} = \hbar \Delta_1 |1\rangle\langle 1| + \hbar \Delta_2 |2\rangle\langle 2| + \frac{\hbar}{2} \Omega_1 (\sigma_{13} e^{-i\vec{k}_1 \cdot \vec{r}} + \sigma_{13}^\dagger e^{i\vec{k}_1 \cdot \vec{r}}) + \frac{\hbar}{2} \Omega_2 (\sigma_{23} e^{-i\vec{k}_2 \cdot \vec{r}} + \sigma_{23}^\dagger e^{i\vec{k}_2 \cdot \vec{r}})$$

$$\dot{\rho}_{AF} = \frac{1}{i\hbar} \left[\tilde{H}_{AF}, \rho_{AF} \right]$$

Trazando sobre grados de libertad atómicos

Operador de Lindblad

$$\mathcal{L}[A]\rho = A\rho A^\dagger - \frac{1}{2}(A^\dagger A \rho + \rho A A^\dagger)$$

$$\dot{\rho}_A = \frac{1}{i\hbar} \left[\tilde{H}, \rho_A \right] + \Gamma_{31} \mathcal{L}[\sigma_{13}] \rho_A$$

$$+ \Gamma_{32} \mathcal{L}[\sigma_{32}] \rho_A$$

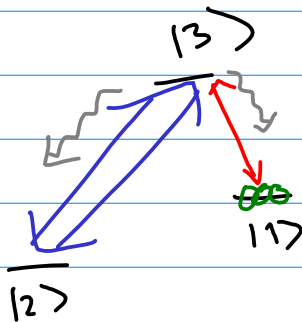
$$+ \gamma_{2dec} \mathcal{L}[\sigma_{22}] \rho_A$$

$$+ \gamma_{1dec} \mathcal{L}[\sigma_{11}] \rho_A$$

$$+ \gamma_{3dec} \mathcal{L}[\sigma_{33}] \rho_A$$

Ecuación maestra.

Caso de láser de prueba débil:



Usando $S_{11} \approx 1$ podemos encontrar la respuesta óptica

Respuesta óptica

$$\langle P^{(+)} \rangle = \underset{\substack{\uparrow \\ \text{densidad} \\ \text{numérica} \\ \text{de átomos}}}{n} \langle J^{(+)} \rangle$$

$$P^{(+)} = \epsilon_0 \chi(\omega) E^{(+)}$$

$$n_{\text{refracción}} = \sqrt{\epsilon_r \mu_r} \approx \sqrt{\epsilon_r} = \sqrt{1 + \chi}$$

